Towards Quantum Control From superposition of programs to quantum recursion

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Motivation

In a Quantum Turing Machine the flow of execution is described by a constant unitary operator.

Developing a notion of quantum control which permits the superposition of finitely many quantum operations.

Building a programming construct which simplifies the presentation of several quantum algorithms, preserving intuition.

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• We are all used to the ever present "if ... then ... else ..." statement, which extracts a boolean value from a predicate and, depending on its truth value, executes one of two statements.

if P(b) then T else Q

• Here b is a bit.

Classical Alternation of Quantum Programs

- The next statement for execution depends on a measurement outcome.
- Construct of probabilistic nature.
- Still resorts to classical information. Therefore not quantum.

measure
$$M[q] = m_1 \rightarrow P_1$$

 $m_2 \rightarrow P_2$
...
 $m_n \rightarrow P_n$
end

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• Here q is a family of qubits.

What if ...

• Recall the definition of classical alternation but with the following changes:

$$egin{array}{l} q \leftarrow |0
angle \ q \leftarrow H[q] \ {f q} {f f} {f q} {f then} \ P {f else} \ Q \end{array}$$

• What can we say about qif and the behaviour of this program?

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Closed Quantum Systems

• Let *H* be a Hilbert space and let $U_0, U_1 : H \to : H$ be unitary operators. Given a qubit **q** define the alternation $Alt_q(U_0, U_1)$ with respect to **q** by

$$\mathit{Alt}_q(\mathit{U}_0,\mathit{U}_1) = \mathsf{\Pi}_0 \otimes \mathit{U}_0 + \mathsf{\Pi}_1 \otimes \mathit{U}_1$$

• It can be represented by a diagonal matrix

$$Alt_q(U_0, U_1) = |0\rangle\langle 0|\otimes U_0 + |1\rangle\langle 1|\otimes U_1 = \begin{pmatrix} U_0 & 0 \ 0 & U_1 \end{pmatrix}$$

• Intuitively, quantum alternation creates a superposition of execution paths of U_0 , U_1 controlled by the basis states of \mathbf{q} .

Examples

• Controlled unitary operations can be written in terms of quantum alternation

qif q_0 then skip else $q_1 \leftarrow U$

• One can also write the toffoli gate as

qif q_0 then skip else qif q_1 then skip else $q_2 \leftarrow N$

• as well as the quantum fourier transform

for i = 1to n do $q_i \leftarrow H$ for k = 2to n - i + 1 do qif q_{k+i-1} then skip else $q_i \leftarrow R_k$

Here R_k is the phase shift gate defined by $R_k = \Pi_0 + e^{i\theta}\Pi_1$ with $\theta = 2\pi/2^k$.

Open Quantum Systems

• States are density operators, and quantum operations are given by superoperators.

- Superoperators represent the most general evolution that an open quantum system can undergo.
- Mathematically superoperators are described by completely positive trace-preserving maps.
- One way to picture superoperators is as the composition of three maps.

$$\rho_{S}(0) \otimes \rho_{B}(0) \xrightarrow{\mathcal{U}} U[\rho_{S}(0) \otimes \rho_{B}(0)]U$$

$$\downarrow^{\text{Tr}_{B}}$$

$$\rho_{S}(0) \xrightarrow{\Phi} \rho_{S}(t)$$

Figure: Superoperator

Kraus representation theorem

• Any superoperator can be written in the form

$$T(\rho) = \sum_{k} E_k \rho E_k^{\dagger}$$

• $E_k = \langle e_k | U | e_0 \rangle$ is an operator on the state space of the state space of the principal system.

• The operators $\{E_k\}$ are known as Kraus decompositions of the quantum operation T. These operators must satisfie the completeness relation

$$\sum_{k} E_{k}^{\dagger} E_{k} \leq I$$

• These decompositions are not unique. One says that two Kraus decompositions are extensionally equal if they represent the same superoperator.

Löwner order on density matrices

• Let D_n be the set of density matrices of dimension n: $D_{\sigma} = \{A \in \mathbb{C}^{n \times n} | A \text{ positive hermitian and tr } A \leq 1\}$

• For matrices $A, B \in \mathbb{C}^{n \times n}$, define $A \sqsubseteq B$ if the matrix B - A is positive. It follows that \sqsubseteq defines a partial order.

• Moreover the poset (D_n, \sqsubseteq) is a complete partial order, i.e, it has least upper bounds of increasing sequences.

• Least upper bounds of increasing sequences coincide with topological limits in the Euclidean topology.

• Any order preserving function on operators will preserve lubs of increasing sequences if it is topologically continuous.

QPL Programming language

Syntax

- high-level features such as loops, recursive procedures, and structured data types.
- functional in nature, statically typed with denotational semantics in terms of a complete partial order of superoperators.

 $\begin{array}{l} QPL \ Terms \ P, Q :::= \mathbf{new \ bit } b := 0 \ | \ \mathbf{new \ qbit } q := 0 \ | \ \mathbf{discard } x \\ | \ b := 0 \ | \ b := 1 \ | \ q_1, ..., q_n * = S \\ | \ \mathbf{skip} \ | \ P; Q \\ | \ \mathbf{if } b \ \mathbf{then } P \ \mathbf{else } Q \ | \mathbf{measure } q \ then \ P \ \mathbf{else } Q \\ | \ \mathbf{while } b \ \mathbf{do } P \\ | \ \mathbf{proc } X : \Gamma \to \Gamma' \{P\} \ \mathbf{in } Q \end{array}$

• A state for a typing context Γ containing n bits and m qubits is given by a $2^n - tuple$ $(A_0, ..., A_{2^n-1})$ of density matrices.

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QPL Programming Language Formal Semantics

• We can assign to QPL a category **Q** which has as its objects tuples of density matrices and whose morphisms are precisely superoperators.

- Identity morphisms are superoperators and, superoperatores are closed under composition.
- $\bullet~\mathbf{Q}$ is equipped with two categorical operations, concatenation and tensor product.
- The Löwner order is naturally extended to matrix tuples. This makes $\mathbf{Q}(\sigma, \sigma')$ into a complete partial order.

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QPL Programming Language

Q atomic morphisms

 \bullet The atomic statements in QPL are interpreted as the following morphisms in ${\bf Q},$

Figure: Morphisms in Q

QPL Programming language

- The recursion can be partially unwound.
- The successive unwindings are given by $F(0), F^2(0), ...$
- Each unwinding is less than the next in the Löwner order, because F is monotone.
- The recursion meaning is given by a least upper bound of the increasing sequence.
- Because density matrices form a complete partial order we are sure that the lubs exist.

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• We can define Quantum Alternation in terms of superoperators as follows :

Given a qubit **q** and two superoperators $T_0, T_1 : S(H) \to S(K)$ then the alternation of T_0 and T_1 with respect to **q** should be the superoperator $A/t_q(T_0, T_1) : S(\mathbf{qbit} \otimes H) \to S(\mathbf{qbit} \otimes K)$,

$$Alt_q(T_0, T_1) :: \rho = \begin{bmatrix} A & B \\ C & D \end{bmatrix} \mapsto \begin{bmatrix} T_0 A & * \\ * & T_1 D \end{bmatrix}$$

 ρ is a state on **qbit** \otimes *H*.

• Quantum alternation should use the information incoded in the classical states of q. If the qubit **q** is in a classical state Π_i with i $\in \{0, 1\}$, the $\{A | t_q(T_0, T_1)\} = I \otimes T_i$. The alternation reduces to a local operation on **qbit** $\otimes S(H)$.

Kraus Semantics

• One general way of defining the semantics of quantum alternation would be in terms of Kraus decompositions.

 \bullet Define a new category ${\bf K}$ with Hilbert spaces as objects and Kraus decompositions as morphisms.

• Composition S o T is defined to be the set obtained from the multiset $X = \{E : F | E \in S, F \in T\}$ by replacing I ocurrences of operator $K \in X$ with \sqrt{IK} .

• Identity is the singleton set containing the usual Identity operator $id_H = \{I\}.$

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• Finally we are in a position to define quantum alternation of two morphisms $S, T : H \to K$ to be the morphism $Alt_q(S, T) : qbit \otimes H \to qbit \otimes K$:

$$Alt_q(S,T) = \left\{ \Pi_0 \otimes \frac{E}{\sqrt{|T|}} + \Pi_1 \otimes \frac{F}{\sqrt{|S|}} | E \in S, F \in T \right\}$$

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QPL and Quantum Alternation

• The semantics of quantum alternation are given by Kraus decompositions.

• The sematics of quantum alternation cannot be lifted to semantics of superoperators because Quantum alternation does not preserve extensional equality.

• Quantum alternation is not compatible with recursion defined in QPL because quantum alternation is not monotone with respect to the Löwner order.

A different approach

Random walks

The simplest random walk is the one-dimensional walk in which a particle moves ona lattice marked by integers \mathbb{Z} , and at each step it moves one position left or right, depending on the flip of a fair coin.

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Quantum Walks

- Hilbert space $H_d \otimes H_p$.
- $H_d = span\{|L\rangle, |D\rangle\}$ denotes the direction space.
- $H_p = span\{|n\rangle : n \in \mathbb{Z}\}$ denotes the position space.
- One step of the walk is represented by the unitary operator $W = T(H \otimes I)$.
- T is a unitary operator in $H_d \otimes H_p$ defined by

$$T|L,n\rangle = |L,n-1\rangle$$
 $T|R,n\rangle = |R,n+1\rangle$

• H is the Hadamard transform in the direction space H_d ,

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$$

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Quantum Walks

Quantum Alternation

- Quantum alternation is disguised in the quantum walk.
- Consider the operators in position space H_p

$$T_L|n\rangle = |n-1\rangle$$
 $T_R|n\rangle = |n+1\rangle$

• The translation operator T can be written as

qif d then $T_L[p]$ else $T_R[p]$

• d represents the direction space and is employed by an external coin system. p represents the position space which we denote as the principal system.

• Furthermore, the single-step walk operator W can be seen as

$$d \leftarrow H[d]$$

qif d then $T_L[p]$ else $T_R[p]$

Quantum Recursion

• We already established that quantum alternation is not compatible with recursion defined in QPL.

• QPL recursion is defined with procedure calls controlled classicaly.

• We may consider this type of recursion as classical recursion of quantum programs.

- Is there a quantum counterpart?
- Consider the program

 $X \leftarrow H[d]$; qif d then $(T_L[p]; X)$ else $(T_R[p]; X)$

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Quantum Recursion

- This program represents a recursive quantum walk.
- The major difference between random and quantum walks is caused by quantum interference.
- Recursive quantum walks exibit a higher-level interference, i.e., interference between infinite paths. Paths containing the recursive walk itself may cancell each other.

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• How do we solve the recursive quantum equation above?

Quantum Recursion

• The number of coin "particles" needed in running a recursive walk is unknown beforehand.

• Need for a method that describes quantum systems with variable number of identical particles.

• Solution: Second Quantization!

• Principle of symmetrisation: States of n identical particles are either completely symmetric or completely antisymmetric with respect to the permutation of the particles.

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• Implementation will depend on the specific choice of coin "particle" : bosons or fermions.

Examples

Quantum while loops

• In classical programming the while-loop

while b do S

can be written as the recursive program

 $X \leftarrow \mathbf{if} \ b \ \mathbf{then} \ X \ \mathbf{else} \ skip$

• Similarly one can write a quantum version

qwhile [c] = |1
angle do U[q]

as

$$X \leftarrow \mathbf{if} \ c \ \mathbf{then} \ skip \ \mathbf{else} \ U[q]; X$$

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• Obviously this two versions are equivalent.



 \bullet A more intersting quantum loop would be

qwhile V[c] = |1
angle do U[q]

as well as

qwhile W[c;q]
$$= |1
angle$$
 do $U[q]$

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