



Hamilton-Jacobi approach to contrast functions and geodetical motion in information geometry

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o	j	f	n	z	g	p	s	c	g	t	b	l	j	z	t	s	l
F	l	o	r	i	o	M	C	i	a	g	l	i	a	c	u	a	m
r	d	j	e	i	x	a	d	s	p	b	i	o	t	n	u	b	q
a	n	g	i	m	u	r	x	t	e	n	s	a	u	g	n	w	r
n	i	w	x	b	l	c	u	w	m	l	P	t	S	z	a	h	m
c	q	j	G	a	i	o	y	n	n	l	a	a	t	v	o	w	x
o	y	o	i	g	s	L	l	y	z	s	t	b	e	m	l	g	t
V	g	w	u	b	l	a	p	F	c	m	r	t	f	q	f	x	g
e	j	r	s	j	h	u	k	a	l	b	i	t	a	z	u	c	i
n	r	w	e	s	t	d	n	b	v	f	z	h	n	t	m	v	b
t	e	f	p	b	F	a	b	i	o	D	i	C	o	s	m	o	m
r	z	z	p	y	n	t	m	o	k	m	a	j	M	d	r	p	w
i	p	a	e	p	w	o	i	M	t	x	V	z	a	c	c	t	d
g	t	v	M	b	x	l	f	e	x	g	i	o	n	l	z	y	o
l	q	k	a	e	w	p	s	l	n	c	t	h	c	v	k	k	m
i	b	n	r	r	t	h	t	e	c	p	a	p	i	d	b	s	h
a	k	s	m	d	r	q	m	j	n	b	l	y	n	b	z	c	i
d	q	D	o	m	e	n	i	c	o	F	e	l	i	c	e	m	a



o	j	f	n	z	g	p	s	c	g	t	b	l	j	z	t	s	l
F	l	o	r	i	o	M	C	i	a	g	l	i	a	c	u	a	m
r	d	j	e	i	x	a	d	s	p	b	i	o	t	n	u	b	q
a	n	g	i	m	u	r	x	t	e	n	s	a	u	g	n	w	r
n	i	w	x	b	l	c	u	w	m	l	P	t	S	z	a	h	m
c	q	j	G	a	i	o	y	n	n	l	a	a	t	v	o	w	x
o	y	o	i	g	s	L	l	y	z	s	t	b	e	m	l	g	t
V	g	w	u	b	l	a	p	F	c	m	r	t	f	q	f	x	g
e	j	r	s	j	h	u	k	a	l	b	i	t	a	z	u	c	i
n	r	w	e	s	t	d	n	b	v	f	z	h	n	t	m	v	b
t	e	f	p	b	F	a	b	i	o	D	i	C	o	s	m	o	m
r	z	z	p	y	n	t	m	o	k	m	a	j	M	d	r	p	w
i	p	a	e	p	w	o	i	M	t	x	V	z	a	c	c	t	d
g	t	v	M	b	x	l	f	e	x	g	i	o	n	l	z	y	o
l	q	k	a	e	w	p	s	l	n	c	t	h	c	v	k	k	m
i	b	n	r	r	t	h	t	e	c	p	a	p	i	d	b	s	h
a	k	s	m	d	r	q	m	j	n	b	l	y	n	b	z	c	i
d	q	D	o	m	e	n	i	c	o	F	e	l	i	c	e	m	a



F l o r i o M C i a g l i a
r a r c P S
n c G o L a F a t r e
c o i u a F a r i a
V e s e d b o D i C o s m o
n t p F a b i o D i C o s m o
r i p t o M e l e V a M
g l e M e l e V i n c i
i a m a r m l n
a D o m e n i c o F e l i c e



- Introduction and motivation
- Generating Potentials for arbitrary tensors
- Hamilton-Jacobi approach to contrast functions



- Pure states are a complex projective space

$$\tilde{h} = \frac{\langle d\psi \otimes d\psi \rangle}{\langle \psi | \psi \rangle} - \frac{\langle \psi | d\psi \rangle}{\langle \psi | \psi \rangle} \otimes \frac{\langle d\psi | \psi \rangle}{\langle \psi | \psi \rangle}$$

$$\pi : \mathcal{H} \rightarrow P\mathcal{H}$$

$$\pi(|\psi\rangle) = \frac{|\psi\rangle\langle\psi|}{\langle\psi|\psi\rangle}$$

- Orthonormal basis $\{|e_j\rangle\}$ $\langle e^j | e_k \rangle = \delta_k^j$

$$z_j = \langle e_j | \psi \rangle \quad \bar{z}^j = \langle \psi | e^j \rangle$$

$$\tilde{h} = \frac{d\bar{z}^j \otimes dz_j}{\|z\|^2} - \frac{(\bar{z}^j dz_j \otimes) \otimes (d\bar{z}_k z^k)}{\|z\|^4}$$



■ non-linear coordinates

$$z_j = \sqrt{p_j} e^{i\varphi_j} \quad \left(p_j \geq 0, \quad \sum_j p_j = 1 \right)$$

$$\tilde{h} = \frac{d\bar{z}^j \otimes dz_j}{\|z\|^2} - \frac{(\bar{z}^j dz_j \otimes) \otimes (d\bar{z}_k z^k)}{\|z\|^4}$$

$$\tilde{h} = \frac{1}{4} \sum_j \frac{1}{p_j} dp_j \otimes dp_j + [dp_j \otimes d\varphi_k + d\varphi_j \otimes d\varphi_k]$$



$$\mathcal{H} \equiv \mathcal{L}^2(\chi) \quad \begin{aligned} |\psi\rangle &\rightarrow \langle x | \psi \rangle = \psi(x) \\ \langle \psi | x \rangle &= \psi^*(x) \end{aligned}$$

- Immersion of a measure space manifold Θ into $P\mathcal{H}$

$$\psi(x|\boldsymbol{\theta}) = \sqrt{p(x|\boldsymbol{\theta})} e^{i\alpha(x|\boldsymbol{\theta})}$$

$$d \ln \psi(x|\boldsymbol{\theta}) = \frac{1}{2} d \ln p(x|\boldsymbol{\theta}) + i d\alpha(x|\boldsymbol{\theta})$$

$$h = g - iw$$

$$g = \frac{1}{4} E_p[(d \ln p) \otimes (d \ln p)] + E_p[d\alpha \otimes d\alpha] - E_p(d\alpha) \otimes E_p(d\alpha)$$

$$w = E_p[d \ln p \wedge d\alpha]$$

$$E_p(df) := \int_{\chi} (df) p dx$$

$$\mathcal{F} = \frac{1}{4} E_p(d \ln p \otimes d \ln p) \quad \text{Fisher-Rao metric}$$



- Statistical Manifold: $\mathcal{M}$
 - Each point $p \in \mathcal{M}$ represents a probability density
- Fisher-Rao metric, a Riemannian metric: g
 - Infinitesimal “distinguishability” of probability densities
- Dualistic Structure: (g, ∇, ∇^*)

$$X, Y, Z \in \mathfrak{X}(\mathcal{M})$$

$$X(g(Y, Z)) = g(\nabla_X Y, Z) + g(Y, \nabla_X^* Z)$$

- Usually defined by perturbation of the hermitean connection

$$\Gamma_{ijk} = {}_g\Gamma_{ijk} \pm T_{ijk}$$



- Statistical manifold: $\mathcal{M}$
- Divergence function: $S : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$
 - $S(\mathbf{x}, \mathbf{y}) \geq 0$
 - $S(\mathbf{x}, \mathbf{y}) = 0 \iff \mathbf{x} = \mathbf{y}$
 - Oriented distance-like function
$$S(x, y) \neq S(y, x)$$
 - If S is differentiable

$$\left. \frac{\partial S}{\partial x^i} \right|_{\mathbf{x}=\mathbf{y}} = \left. \frac{\partial S}{\partial y^i} \right|_{\mathbf{x}=\mathbf{y}} = 0$$

- S is a generating function
- Generates tensors intrinsically $\frac{\partial^2 S}{\partial x^i \partial x^j} = g_{ij}$



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- Some examples
- Hessian Riemannian Structures^[1]:
 - $\mathbf{x} = \{x^i\}$ Local coordinates on $U \subset \mathcal{M}$
 - Convex function: $f : U \rightarrow \mathbb{R}$

$$g_{ij}(\mathbf{x}) := \frac{\partial^2 f}{\partial x^i \partial x^j}$$

This is not a tensor!

- Kähler Potential

$$\omega = \frac{\partial^2 f}{\partial z \partial \bar{z}}$$

Quote from Wikipedia: “There is no comparable way of describing a general Riemannian metric in terms of a single function.”

[1]: J.J. Duistermaat. On Hessian Riemannian Structures. Asian J. Math, 5 (1) (2001), pp. 79-92



- Two copies of the manifold to define a two-point function

$$S : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$$

- Provides several canonical structures:

$$X \in \mathfrak{X}(\mathcal{M})$$

$$\pi_L : \mathcal{M} \times \mathcal{M} \rightarrow \mathcal{M}$$

$$\text{Left-lift: } X_L \in \mathfrak{X}(\mathcal{M} \times \mathcal{M})$$

$$\pi_R : \mathcal{M} \times \mathcal{M} \rightarrow \mathcal{M}$$

$$\text{Right-lift: } X_R \in \mathfrak{X}(\mathcal{M} \times \mathcal{M})$$

Diagonal embedding:

$$d : \mathcal{M} \hookrightarrow \mathcal{M} \times \mathcal{M}$$

$$p \mapsto (p, p)$$

Restriction to the diagonal:

$$S|_d := d^* S$$



- Consider the contrast function $S : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$

$$(\mathcal{L}_{X_L} \mathcal{L}_{Y_L} S) \big|_d := g(X, Y)$$

- It is clearly f-linear in X
- Let us show that it is also f-linear in Y

$$\begin{aligned} \mathcal{L}_{X_L} \mathcal{L}_{(fY)_L} S &= \mathcal{L}_{X_L} (f \mathcal{L}_{Y_L} S) \\ &= f \mathcal{L}_{X_L} \mathcal{L}_{Y_L} S + \mathcal{L}_{X_L} f \cdot \mathcal{L}_{Y_L} S \end{aligned}$$



- Consider the contrast function $S : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$

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$$\begin{aligned} (\mathcal{L}_{X_L} \mathcal{L}_{(fY)_L} S) \big|_d &= (\mathcal{L}_{X_L} (f \mathcal{L}_{Y_L} S)) \big|_d \\ &= (f \mathcal{L}_{X_L} \mathcal{L}_{Y_L} S) \big|_d + (\mathcal{L}_{X_L} f \cdot \cancel{\mathcal{L}_{Y_L} S}) \big|_d \\ &= fg(X, Y) \end{aligned}$$

Note: In the original image, a red arrow points from the '0' in the second term of the second line to the '0' in the third line, indicating that the term is zero.



- Consider the contrast function $S : \mathcal{M} \times \mathcal{M} \rightarrow \mathbb{R}$

$$(\mathcal{L}_{X_L} \mathcal{L}_{Y_L} S)|_d := g(X, Y)$$

- It is clearly f-linear in X
- Let us show that it is also f-linear in Y

$$\begin{aligned} (\mathcal{L}_{X_L} \mathcal{L}_{(fY)_L} S)|_d &= (\mathcal{L}_{X_L} (f \mathcal{L}_{Y_L} S))|_d \\ &= (f \mathcal{L}_{X_L} \mathcal{L}_{Y_L} S)|_d + (\mathcal{L}_{X_L} f \cdot \cancel{\mathcal{L}_{Y_L} S})|_d \quad \begin{matrix} 0 \\ \nearrow \end{matrix} \\ &= fg(X, Y) \end{aligned}$$

- One can define more rank 2 tensors:

$$(\mathcal{L}_{X_R} \mathcal{L}_{Y_R} S)|_d = (\mathcal{L}_{X_L} \mathcal{L}_{Y_L} S)|_d = -(\mathcal{L}_{X_R} \mathcal{L}_{Y_L} S)|_d = g(X, Y)$$



- This procedure can be used for tensors of order 3
- One needs to find suitable linear combinations to cancel the non-tensorial terms

$$(\mathcal{L}_{X_L} \mathcal{L}_{Y_R} \mathcal{L}_{Z_R} S - \mathcal{L}_{X_R} \mathcal{L}_{Y_L} \mathcal{L}_{Z_L} S)|_d = T(X, Y, Z)$$

or

$$(\mathcal{L}_{X_L} \mathcal{L}_{Y_L} \mathcal{L}_{Z_R} S - \mathcal{L}_{X_R} \mathcal{L}_{Y_R} \mathcal{L}_{Z_L} S)|_d = T(X, Y, Z)$$

- There are 8 possible combinations that lead to tensorial quantities.
 - They define the same symmetric tensor up to a sign
 - In information geometry this is used to define pairs of dual connections



- Can one go to higher order tensors?
 - For instance the Riemann tensor

- Can one generate non-symmetric tensors?



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- Find a canonical contrast function
 - Given g, T
 - Construct a two point function with vanishing derivatives

- Hamilton's Principal function

$$S[\gamma] = \int_{t_{\text{in}}}^{t_{\text{fin}}} \mathfrak{L}(\gamma(t), \dot{\gamma}(t)) dt$$

- Solution of the equations of motion γ
- $\gamma(0) = \mathbf{x} \quad \gamma(1) = \mathbf{y}$

$$S[\mathbf{x}, \mathbf{y}] = \int_0^1 \mathfrak{L}(\gamma(t), \dot{\gamma}(t)) dt$$

- $dS = p_i(\mathbf{x}, \mathbf{y}) dx^i - P_i(\mathbf{x}, \mathbf{y}) dy^i$

$$\mathbf{x} = \mathbf{y} \quad \gamma(t) = \mathbf{x} \quad \Rightarrow \quad \begin{aligned} p_i(\mathbf{x}, \mathbf{y})|_{\mathbf{x}=\mathbf{y}} &= 0 \\ P_i(\mathbf{x}, \mathbf{y})|_{\mathbf{x}=\mathbf{y}} &= 0 \end{aligned}$$



■ Choose the Lagrangian

$$\mathfrak{L}(x, v) = \frac{1}{2}g_{ij}v^i v^j + \frac{1}{6}T_{ijk}v^i v^j v^k$$

■ Compute derivatives

$$\frac{\partial S}{\partial y^j} = g_{jk}(\mathbf{y})v_{\text{fin}}^k + \frac{1}{2}T_{jkl}(\mathbf{y})v_{\text{fin}}^k v_{\text{fin}}^l$$

- We need the relation between $\mathbf{v}_{\text{fin}}$ and $\mathbf{x}, \mathbf{y}$

- Taylor expansion

$$\gamma(t) = \gamma(1) + \mathbf{v}_{\text{fin}}(1-t) + \frac{1}{2}\dot{\mathbf{v}}_{\text{fin}}(1-t)^2 + \frac{1}{6}\ddot{\mathbf{v}}_{\text{fin}}(1-t)^3$$

$$\Rightarrow \mathbf{v}_{\text{fin}} = \mathbf{y} - \mathbf{x} + \frac{1}{2}\dot{\mathbf{v}}_{\text{fin}} - \frac{1}{6}\ddot{\mathbf{v}}_{\text{fin}}$$

- Eq. of motion

$$\dot{v}^l = -T_{jk}^l v^k \dot{v}^j - \Gamma_{jk}^l v^j v^k - \frac{1}{6}g^{lr}A_{rjks}v^j v^k v^s$$



- Assuming that $\dot{\mathbf{v}}$ is analytic in $\mathbf{v}$

$$\Delta^k := y^k - x^k$$

$$\begin{aligned} \frac{\partial S}{\partial y^j} = & g(\mathbf{y})_{jk} \Delta^k - \frac{1}{2} \Gamma_{jkl}(\mathbf{y}) \Delta^k \Delta^l + \frac{1}{2} T_{jkl}(\mathbf{y}) \Delta^k \Delta^l \\ & + \frac{1}{6} \Gamma_{jkr}(\mathbf{y}) \Gamma_{ls}^r(\mathbf{y}) \Delta^k \Delta^l \Delta^s - \frac{1}{12} A_{jkl s}(\mathbf{y}) \Delta^k \Delta^l \Delta^s \end{aligned}$$

$$\left. \frac{\partial^2 S}{\partial x^i \partial y^j} \right|_d = -g_{ij}(\mathbf{x})$$

$$\left. \frac{\partial^3 S}{\partial x^l \partial x^k \partial y^j} \right|_d = -\Gamma_{jkl} + T_{jkl}$$

By similar calculations:

$$\left. \frac{\partial^3 S}{\partial y^l \partial y^k \partial x^j} \right|_d = -\Gamma_{jkl} - T_{jkl}$$



- Find a combination of 4th order derivatives that is f-linear

$$\mathcal{L}_{X_L} \mathcal{L}_{Y_R} \mathcal{L}_{Z_L} \mathcal{L}_{W_R} S|_d =: LRLR$$

- Two possible linear combinations:

$$Q_1(X, Y, Z, W) = (LRRL - RLLR) + (LRRR - RLLL) \\ + (LRLL - RLRR) + (LRLR - RLRL)$$

$$Q_2(X, Y, Z, W) = (LLLL - RRRR) + (LLLR - RRRL) \\ + (LLRL - RRRL) + (LLRR - RRLL)$$



- Find a combination of 4th order derivatives that is f-linear

$$\mathcal{L}_{X_L} \mathcal{L}_{Y_R} \mathcal{L}_{Z_L} \mathcal{L}_{W_R} S|_d =: LRLR$$

- Two possible linear combinations:

$$\begin{aligned} Q_1(X, Y, Z, W) = & (LRRL - RLLR) + (LRRR - RLLL) \\ & + (LRLL - RLRR) + (LRLR - RLRL) \end{aligned}$$

$$\begin{aligned} Q_2(X, Y, Z, W) = & (LLLL - RRRR) + (LLLR - RRRL) \\ & + (LLRL - RRLR) + (LLRR - RRLL) \end{aligned}$$

- After a lengthy calculation:

$$Q_1 = Q_2 = 0$$



- Choose a lagrangian

$$\mathfrak{L}(x, v) = \frac{1}{2}g_{ij}v^iv^j + \frac{1}{6}T_{ijk}v^iv^jv^j + \frac{1}{24}C_{ijkl}v^iv^jv^jv^l$$



- Choose a lagrangian

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- After a lengthy calculation:

$$Q_1 = Q_2 = 0$$

- What about non-symmetric tensors?



- F.M. Ciaglia, F. di Cosmo, D. Felice, S. Mancini, G. Marmo and J.M. Pérez-Pardo. *Hamilton-Jacobi approach to Potential Functions in Information Geometry*. **J. Math. Phys.** **58**, 063506. (2017).
- P. Facchi, R. Kulkarni, V.I. Man'ko, G. Marmo, E.C.G. Sudarshan and F. Ventriglia. *Classical and quantum Fisher information in the geometrical formulation of quantum mechanics*. **Phys. Lett. A** **374**, (2010), 4801-4803