

# Around the dynamics of Painlevé's equations

Helena Reis

(joint work with A. Belloto and J. Rebelo - Univ. Toulouse)

September 7, 2017

## Painlevé (classical) equations

$$(P_I) \quad y'' = 6y^2 + x$$

$$(P_{II}) \quad y'' = 2y^3 + xy + \alpha$$

$$(P_{III}) \quad y'' = \frac{1}{y}(y')^2 - \frac{1}{x}y' + \frac{1}{x}(\alpha y^2 + \beta) + \gamma y^3 + \frac{\delta}{y}$$

$$(P_{IV}) \quad y'' = \frac{1}{2y}(y')^2 + \frac{3}{2}y^3 + 4xy^2 + 2(x^3 - \alpha)y + \frac{\beta}{y}$$

$$(P_V) \quad y'' = \left(\frac{1}{2y} + \frac{1}{y-1}\right)(y')^2 - \frac{1}{x}y' + \frac{(y-1)^2}{x^2} \left(\alpha y + \frac{\beta}{y}\right) \frac{\gamma y}{x} + \frac{\delta y(y+1)}{y-1}$$

$$(P_{VI}) \quad y'' = \frac{1}{2} \left(\frac{1}{y} + \frac{1}{y-1} + \frac{1}{y-x}\right)(y')^2 - \left(\frac{1}{x} + \frac{1}{x-1} + \frac{1}{y-x}\right)y' \\ + \frac{y(y-1)(y-x)}{x^2(x-1)^2} \left(\alpha + \frac{\beta x}{y^2} + \frac{\gamma(x-1)}{(y-1)^2} + \frac{\delta x(x-1)}{(y-x)^2}\right)$$

$$\alpha, \beta, \gamma, \delta \in \mathbb{C}$$

# Painlevé (classical) equations

**Purpose:** To work out the dynamics associated with these equations.

## **Applications:**

- (1) to determine when the Galois-Malgrange pseudogroup is maximal
- (2) to obtain accurate asymptotic estimates for solutions

# Galois-Malgrange pseudogroup

**Galois-Malgrange pseudogroup:** *generalization of the differential Galois group associated with linear equations (algebraic definition whose dynamical meaning is not obvious)*

- *(Virtually) non-solvable differential Galois group  $\Rightarrow$  the solutions cannot be “integrated by quadratures”*
- *Painlevé equations are not linear: problem of irreducibility of Painlevé “transcendents”*

# Irreducibility of Painlevé transcendents

- Umemura notion of irreducibility  $\Rightarrow$  Umemura results + works by the Japanese school
- Galois-Malgrange pseudogroup
  - (1) general definition implying all other forms of irreducibility
  - (2) more directly related with dynamics in that it is defined as a “closure” for holonomy groups in the spirit of dynamical interpretations of Galois theory
  - (3) coincides with the Galois group if the equation is linear (Malgrange theorem, not trivial)
  - (4) it can also be viewed as a “measure” - albeit a coarse one - of the dynamical complexity of the equations

# Galois-Malgrange pseudogroup

*How to compute Galois-Malgrange pseudogroup?*

**For order 2 equations, we have**

Theorem (Casale)

*The Galois-Malgrange pseudogroup is maximal if and only if*

- (1) there is no algebraic codimension 1 (multi)foliation invariant by the equation and*
- (2) there is no algebraic transverse (multi)affine structure invariant by the equation*

## Galois-Malgrange pseudogroup

Consider the Painlevé equation  $P_I$

$$y'' = 6y^2 + x.$$

This equation is equivalent to the vector field

$$X_I = \frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + (6y^2 + x) \frac{\partial}{\partial z}$$

in  $\mathbb{C}^3$ .

*The Galois-Malgrange pseudogroup is maximal if and only if*

- (1) there is no algebraic codimension 1 (multi)foliation whose leaves contain the orbits of  $X_I$*
- (2) there is no algebraic transverse (multi)affine structure for  $\mathcal{F}_I$ , the foliation associated to  $P_I$*

# Galois-Malgrange pseudogroup

*Toy-model case:* **Airy equation**

$$y'' = xy$$

*whose associated vector field is*

$$X_A = \frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + xy \frac{\partial}{\partial z}$$



# Galois-Malgrange pseudogroup

## Some known results:

- **Airy equation:** the Galois-Malgrange pseudogroup is maximal (“classical” result with algebraic computations)
- $P_I$  **equation:** the Galois-Malgrange pseudogroup is maximal (Casale)
- $P_{II}$ - $P_{VI}$  **equation:** the Galois-Malgrange pseudogroup is maximal for “generic values” of the parameters (Casale, Roques, etc)
- $P_{VI}$  **equation:** complete characterization by Cantat-Loray

# Problems about “asymptotic estimates”

## $P_I, P_{II}, P_{IV}$ and modified $P_{III}, P_V$ equations

- they have meromorphic solutions defined on  $\mathbb{C}$  (i.e. the solutions are holomorphic functions  $\varphi : \mathbb{C} \rightarrow \mathbb{CP}(1)$ )
- they have an essential singularity at “ $T = \infty$ ”



- \* study of their asymptotic behavior (Nevlinna theory)
- \* purely dynamical problem: ergodic properties

## Remark

The above mentioned questions will not be posed for  $P_{VI}$  - solutions do not have a maximal domain of definition on  $\mathbb{C}$ .

# Airy Equation

Vector field in correspondence to the Airy equation

$$X_A = \frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + xy \frac{\partial}{\partial z}$$

Alternative (more dynamical-geometric proof) proof of the maximality of the Galois-Malgrange pseudogroup (coinciding with the differential Galois group).

# Airy Equation

## Theorem

Consider the (meromorphic) extension of  $X_A$  to  $\mathbb{CP}(3) = \mathbb{C}^3 \cup \Delta_\infty$ . There are two points  $p, q \in \Delta_\infty$  such that the following holds: given two neighborhoods  $V_p, V_q$  of  $p, q$  on  $\mathbb{CP}(3)$ , respectively, and an integral curve  $\phi$  for  $X_A$ , we have

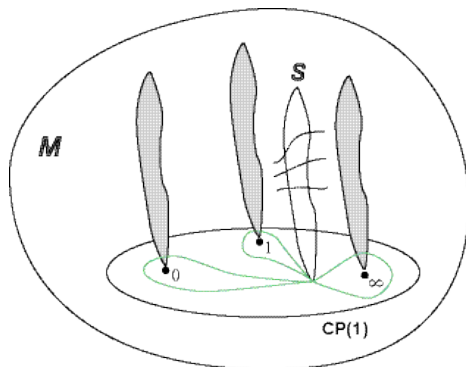
$$\lim_{r \rightarrow \infty} \frac{\text{Area}(\{T \in B(r) : \phi(T) \in V_p \cup V_q\})}{\text{Area}(B(r))} = 1$$

## Remark

$X_A$  is “very non-ergodic” - its solutions are “confined” in a probabilistic sense.

# Global properties of Painlevé equations

Riemann-Hilbert picture (Okamoto compactification) for  $P_{VI}$



# Global properties of Painlevé equations

**Holonomy group:** *generated by two independent holonomy (algebraic) maps  $f = h_{\gamma_0}$  and  $g = h_{\gamma_\infty}$*

$$f : S \rightarrow S$$

$$g : S \rightarrow S$$

*Let  $\bar{S}$  stands for the compactification of  $S$ . Then  $f, g : \bar{S} \rightarrow \bar{S}$  are **birational maps**.*

## Cantat-Loray viewpoint

### Dynamical study if the action of $\Gamma = \langle f, g \rangle$ on $\bar{S}$

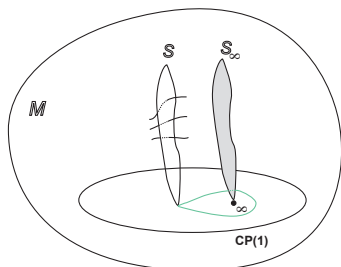
- **If** there exists codimension 1 foliation invariant by  $\mathcal{F}_{P_{VI}}$ , **then** there exists a 1-dimensional foliation on  $\bar{S}$  invariant by  $\Gamma$ .
- **If** there exists a transverse affine structure for  $X_{P_{VI}}$ , **then** there exists a (singular) affine structure on  $\bar{S}$  invariant by  $\Gamma$

*By studying the dynamics of  $\Gamma$  Cantat and Loray rule out the existence of these objects for all parameters but the case corresponding to Picard solutions.*

*There is no codimension 1 foliation invariant by  $X_{P_{VI}}$ , not even among foliations with entire coefficients.*

# Global properties of Painlevé equations

Riemann-Hilbert picture (Okamoto compactification) for  $P_I/P_{II}$



*base simply connected*  $\Rightarrow$  *the holonomy maps are trivial*  
 $\Rightarrow$  *plenty of invariant **holomorphic** foliations /*  
*transverse affine structures*



# Our approach

**Problem:** How to detect if, say an invariant codimension 1 foliation, is algebraic?

- Notion of dynamics at infinity
- incidentally this “dynamics” at infinity is exactly the object that controls the asymptotic behaviour of solutions
- add a fiber over the “missing points” (if they were missing) and study the foliation in the neighbourhood of this fiber

## The case of Airy equation

The Airy equation defined a 1-dimensional foliation  $\mathcal{F}$  on  $\mathbb{C}^3$ , induced by

$$X_A = \frac{\partial}{\partial x} + z \frac{\partial}{\partial y} + xy \frac{\partial}{\partial z}$$

In standard coordinates for  $\mathbb{C}\mathbb{P}(3)$  around  $\Delta_\infty$

$$x_1 = \frac{1}{x}, \quad y_1 = \frac{y}{x}, \quad z_1 = \frac{z}{x}$$

where  $\Delta_\infty \subseteq \{x_1 = 0\}$ , we have

$$X_{A,1} = \frac{1}{x_1} \left[ -x_1^3 \frac{\partial}{\partial x_1} + (x_1 z_1 + y_1 x_1^2) \frac{\partial}{\partial y_1} + (y_1 - x_1^2 z_1) \frac{\partial}{\partial z_1} \right].$$

## The case of Airy equation

$\mathcal{F}$  extends to a (singular) foliation on  $\mathbb{C}^3 \cup \Delta_\infty \simeq \mathbb{CP}(3)$

$$X_{A,1}|_{\Delta_\infty} = y_1 \frac{\partial}{\partial z}$$

Some issues on  $\mathcal{F}$ :

- global dynamics on  $\Delta_\infty$  is “simple”
- $\mathcal{F}$  has singular points that may “conceal” non-trivial dynamics
- $\mathcal{F}$  possesses compact leaves on  $\Delta_\infty$  which may carry holonomy
- there is a foliated affine structure on the leaves of  $\mathcal{F}$

## Singularities:

In coordinates  $(x_1, y_1, z_1)$  for  $\mathbb{CP}(3)$ , where  $\{x_1 = 0\} \subseteq \Delta_\infty$

$$\text{Foliation: } X_{A,1} = -x_1^3 \frac{\partial}{\partial x_1} + (x_1 z_1 + y_1 x_1^2) \frac{\partial}{\partial y_1} + (y_1 - x_1^2 z_1) \frac{\partial}{\partial z_1}.$$

1.  $\{x_1 = 0, y_1 = 0\} = C_1 \rightarrow$  curve of singularities
2. nilpotent singularities

In coordinates  $(x_2, y_2, z_2)$  for  $\mathbb{CP}(3)$ , where  $\{x_2 = 0\} \subseteq \Delta_\infty$

$$\text{Foliation: } X_{A,2} = (y_2^2 - x_2 y_2 z_2) \frac{\partial}{\partial x_2} - y_2^2 z_2 \frac{\partial}{\partial y_2} + (x_2 - y_2 z_2^2) \frac{\partial}{\partial z_2}.$$

1.  $\{x_2 = 0, y_2 = 0\} = C_2 \rightarrow$  curve of singularities
2. plenty of degenerate singularities
3. reduction of singularities

## Resolution of singularities:

1. 2 ramified blow-ups along  $C_1$  and  $C_2$  -  $\Delta_\infty \cup \Delta_1 \cup \Delta_2$  total divisor
2. the resulting foliation possesses only 3 singular points:  $q, p_1, p_2$
3.  $q \in \Delta_\infty \cap \Delta_1 \cap \Delta_2$  - simple singular point
4.  $p_1 \in \Delta_1 \cap \Delta_2$ ;  $p_2 \in \Delta_2 \setminus (\Delta_\infty \cup \Delta_1)$  - saddle-node singularities
5.  $L_1 = \Delta_1 \cap \Delta_2$  - rational invariant curve
6. there exists  $L_2 \subseteq \Delta_2$  containing  $p_1$  and  $p_2$  - invariant rational curve

*The invariant rational curves  $L_1$  and  $L_2$  form a sort of “attractor” of the dynamics (every leaf intersects a neighbourhood of this curve)*

## Fundamental dynamical issue:

*To understand the dynamics of the foliation on a neighborhood of  $L_1 \cup L_2$ .*

- *one non-trivial holonomy map  $h : (\mathbb{C}^2, 0) \rightarrow (\mathbb{C}^2, 0)$*
- *highly resonant saddle-nodes (eigenvalues:  $0, 1, -1$ )*